



# Fluorination of Graphene, Gas Sensing of Carbon-based Materials, and Spectral Analysis of Water/Oxides Surface: Physics of Adsorption and Interfaces

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Department of Physics

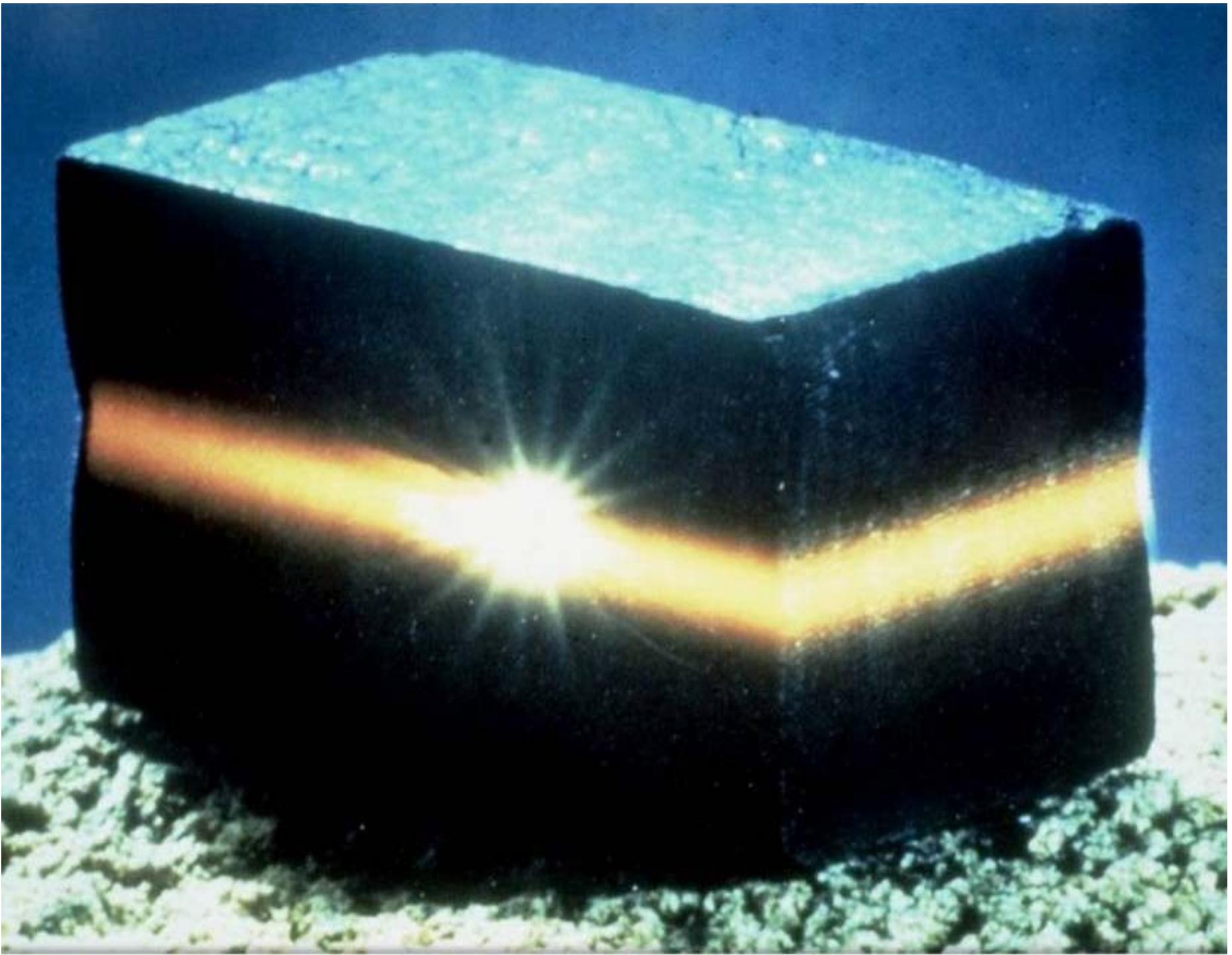
and

Department of Materials Science and Engineering

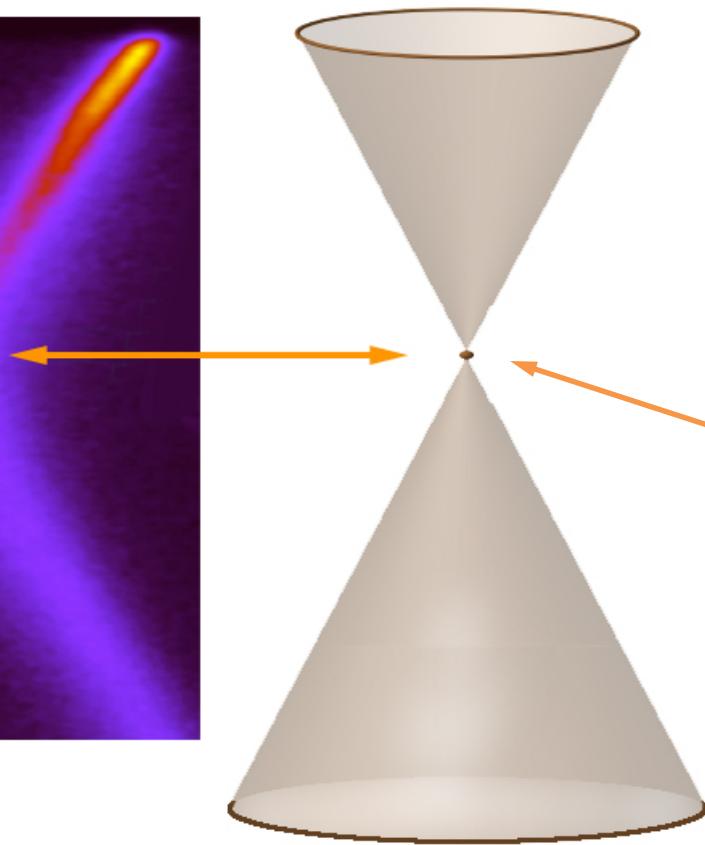
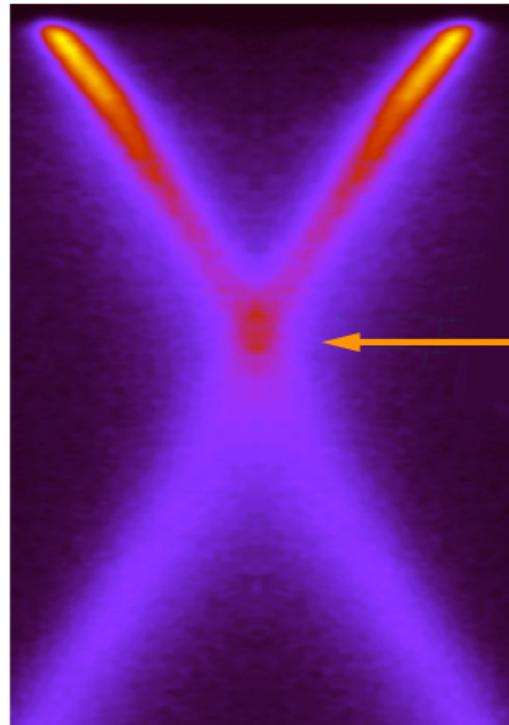
and

Materials Research Institute

Penn State

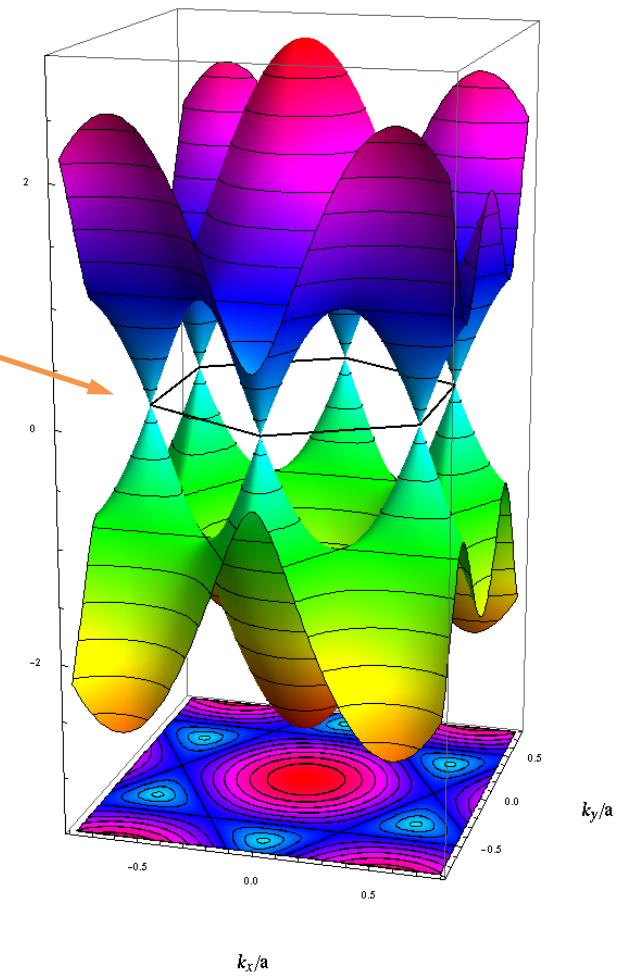


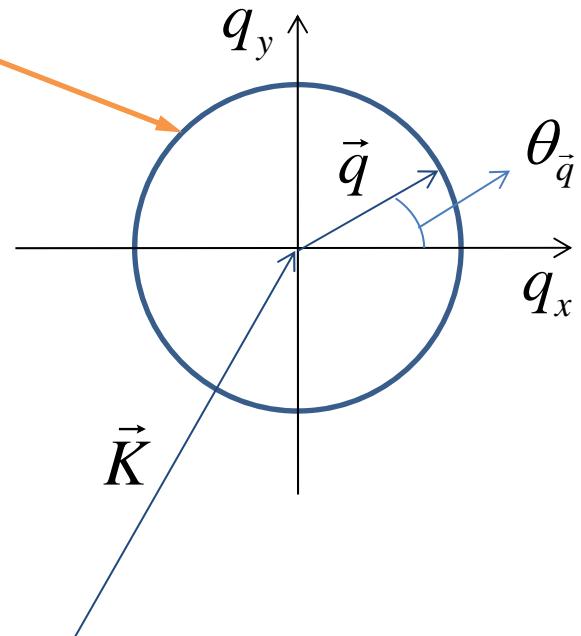
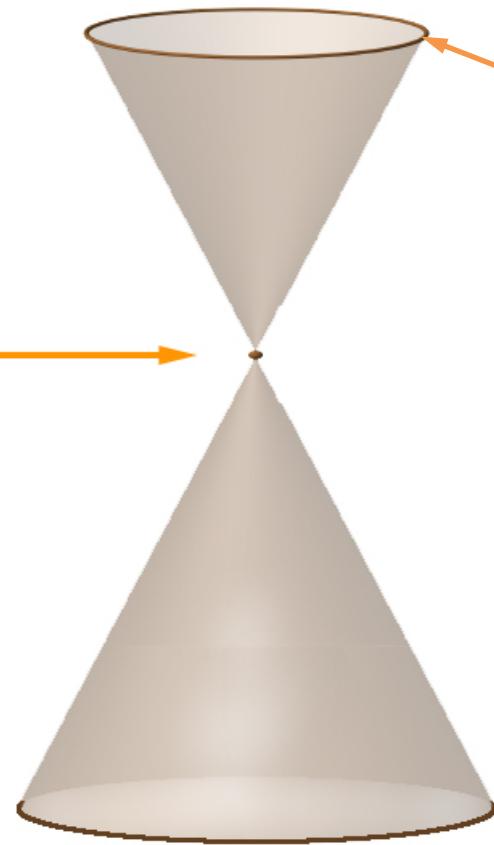
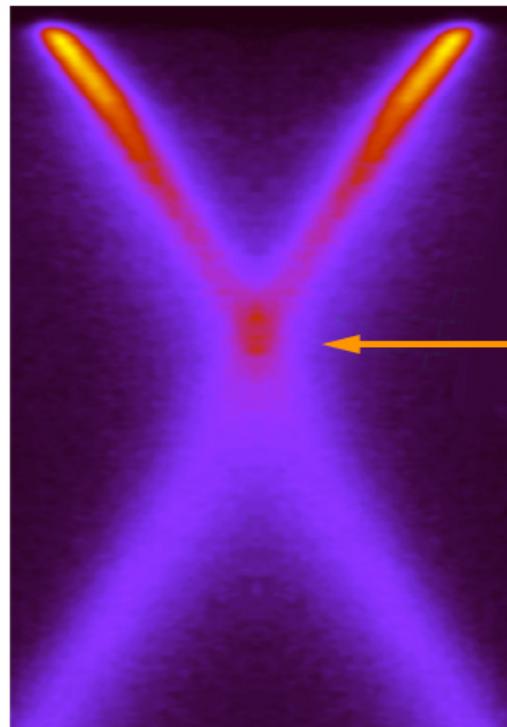
Courtesy: Richard Kaner, UCLA



ARPES Measurement  
 (Aaron Bostwick and Eli Rotenberg, LBNL)

$$\gamma_0 \approx 3 \text{ eV}$$





# Massless Dirac Fermions

$$H_{\vec{k}} = \begin{bmatrix} 0 & -\gamma_0 \phi(\vec{k}) \\ -\gamma_0 \phi^*(\vec{k}) & 0 \end{bmatrix} \quad \vec{k} = \vec{K} + \vec{q}$$

$|\vec{K}| \gg |\vec{q}|$

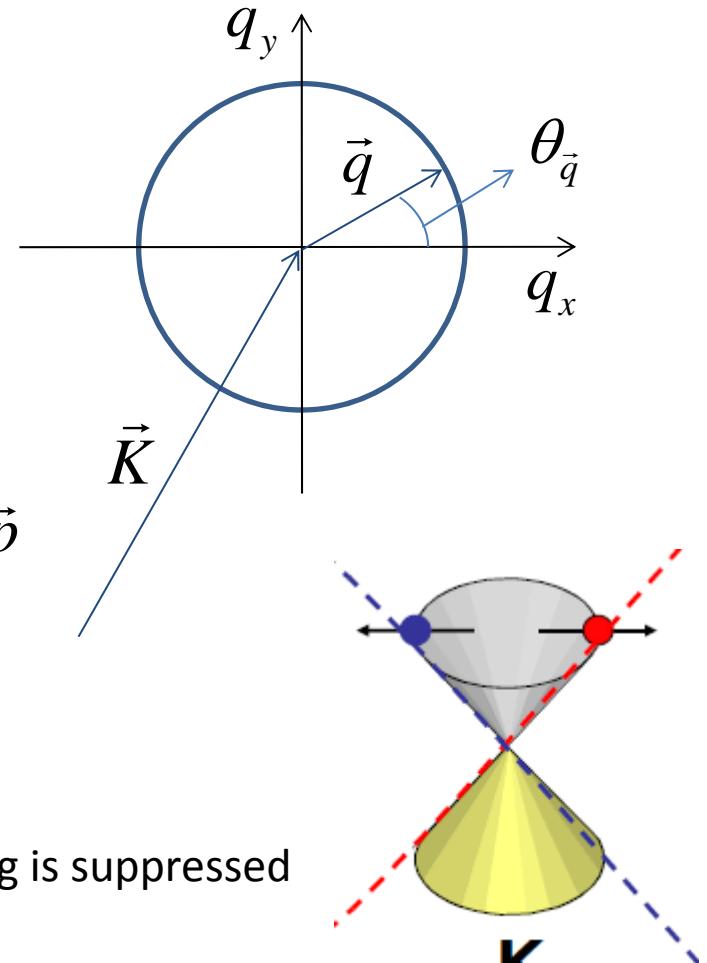
$$H_{\vec{q}} = \begin{bmatrix} 0 & -\frac{3\gamma_0 a}{2} |\vec{q}| e^{i\theta_{\vec{q}}} \\ -\frac{3\gamma_0 a}{2} |\vec{q}| e^{-i\theta_{\vec{q}}} & 0 \end{bmatrix} = v_F \vec{\sigma} \cdot \vec{p}$$

$$\varepsilon_{\pm}(\vec{q}) = \pm \frac{3\gamma_0 a}{2} |\vec{q}| = \pm \hbar v_F |\vec{q}|$$

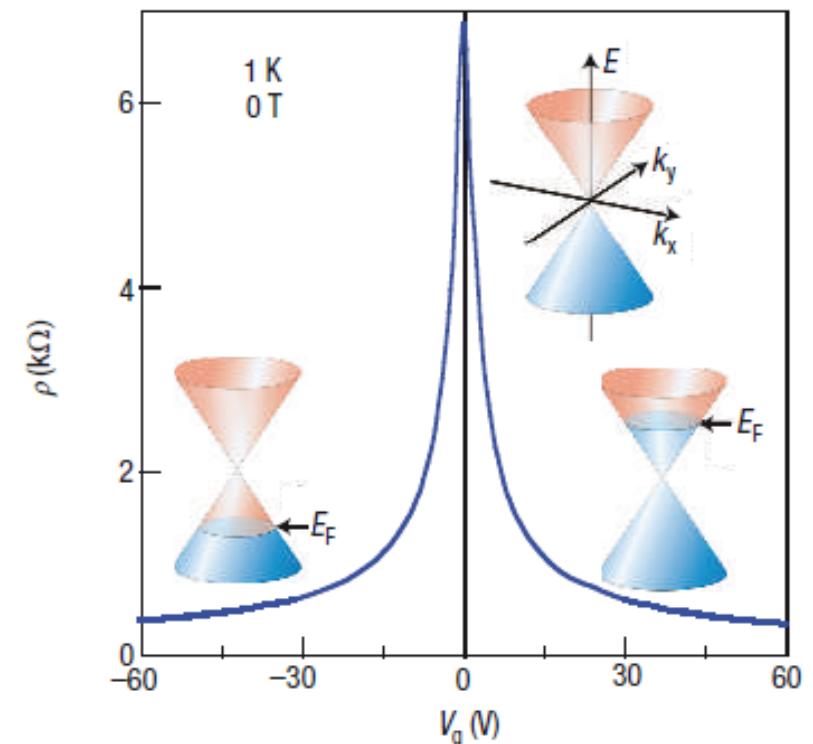
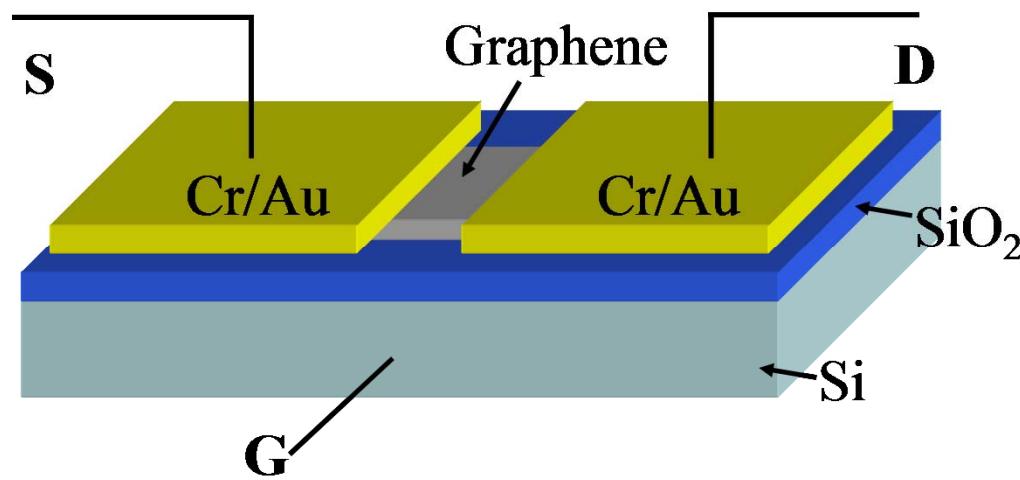
$$|\pm, \vec{q}\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\frac{\theta_{\vec{q}}}{2}} \\ \pm e^{i\frac{\theta_{\vec{q}}}{2}} \end{pmatrix}$$

$$\langle +, \vec{q} | | +, -\vec{q} \rangle = 0 \quad !!$$

Transport:  
Backscattering is suppressed

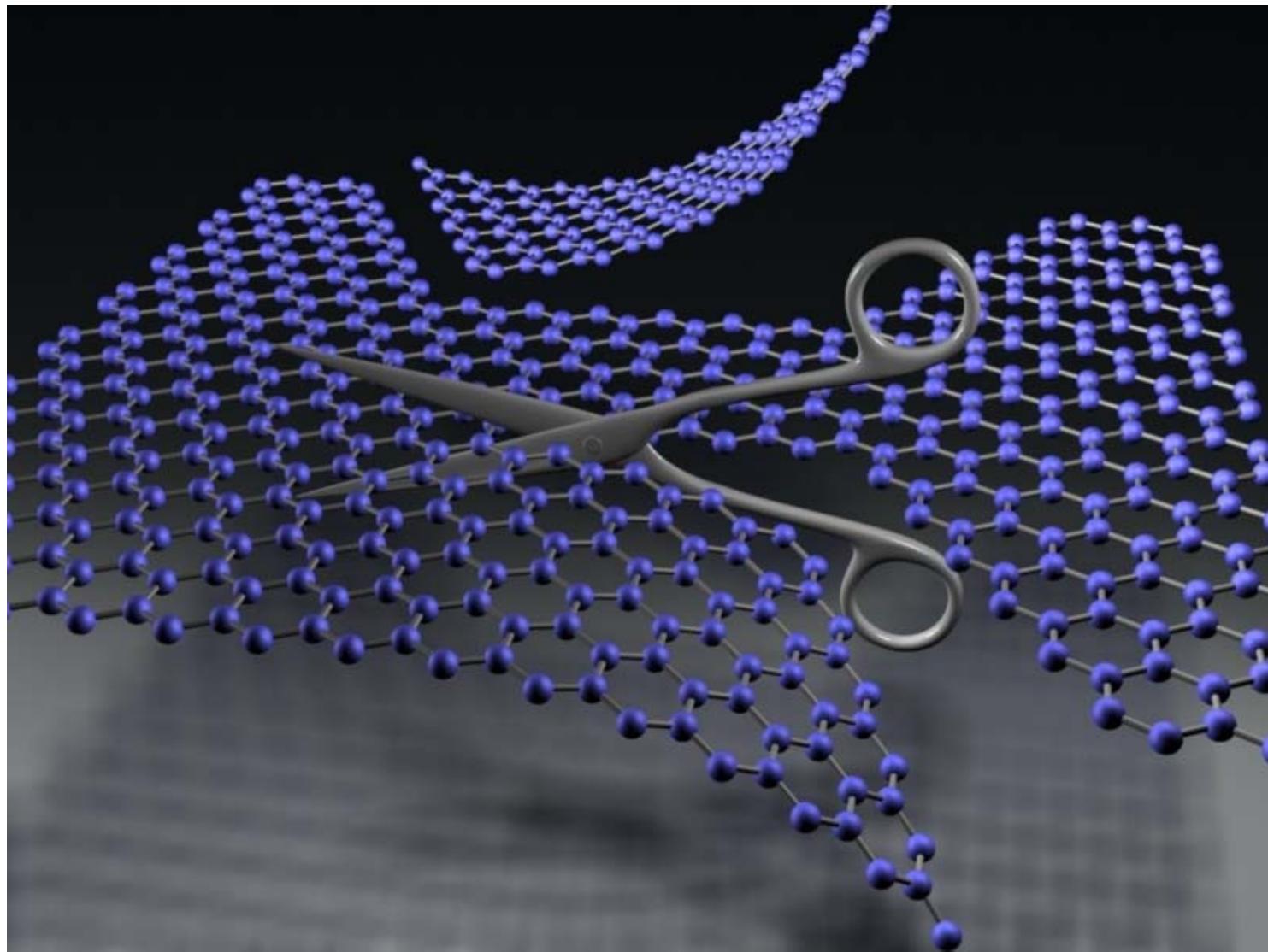


# Carrier density controlled by gate voltage



Hugo Romero, et al., *ACS Nano* (September 30, 2008).

A. Geim and K. Novoselov.  
“The rise of graphene.”  
*Nat Mater* 6, 183 (2007)



Graphics courtesy of  
Kostya Novoselov

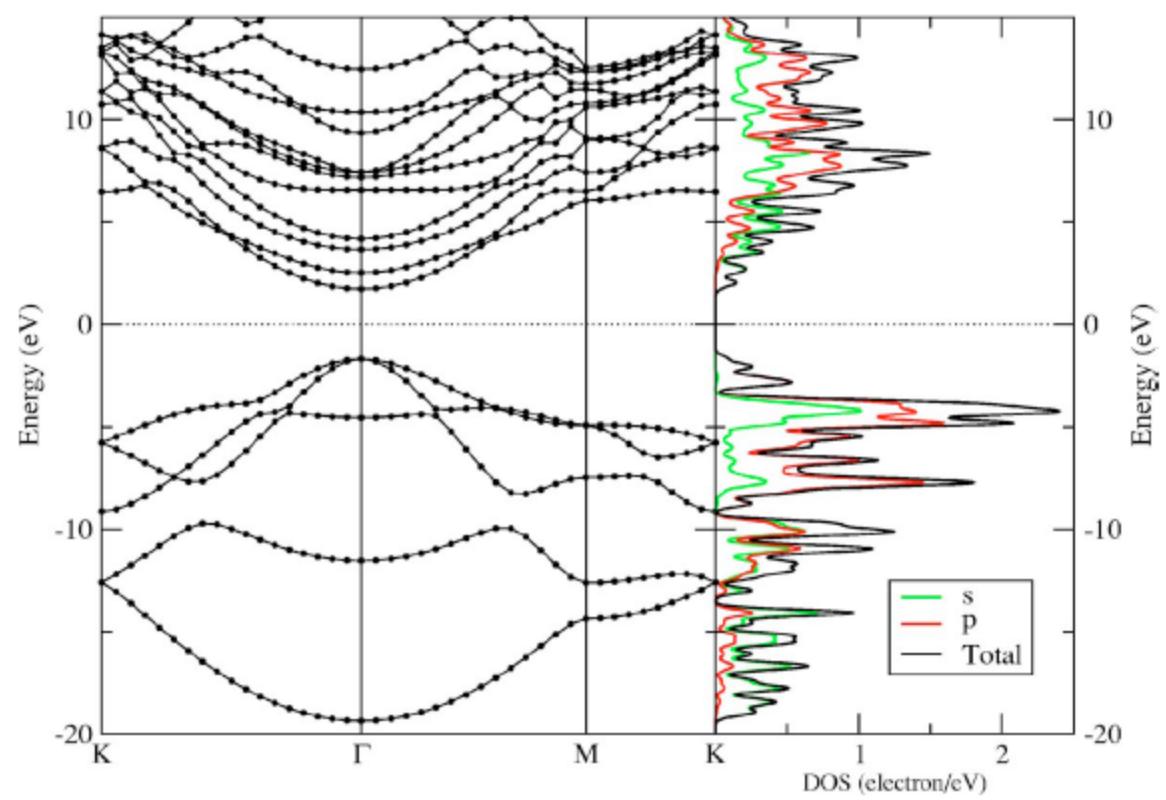
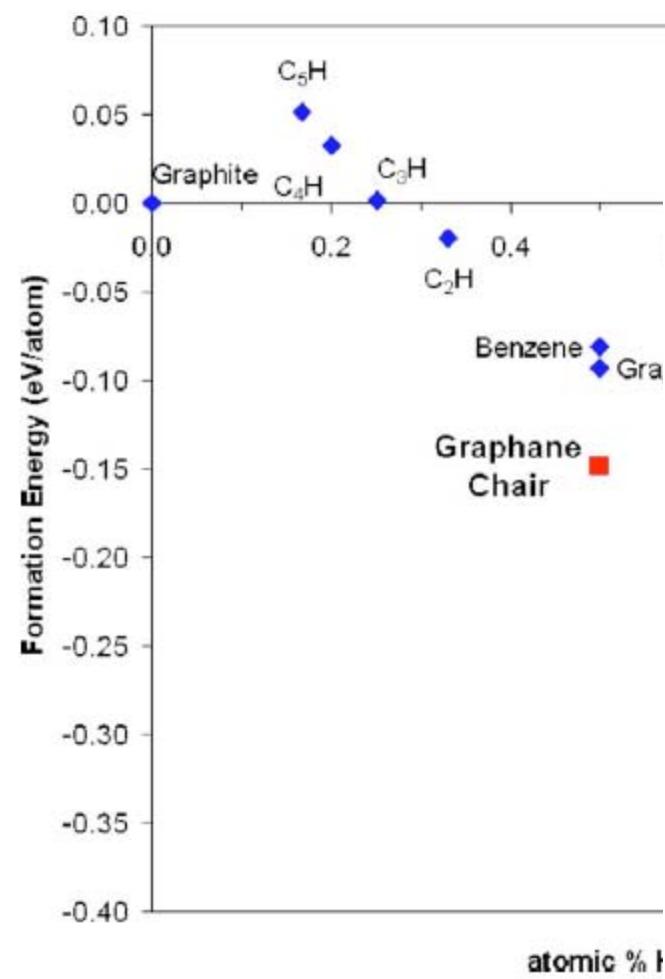
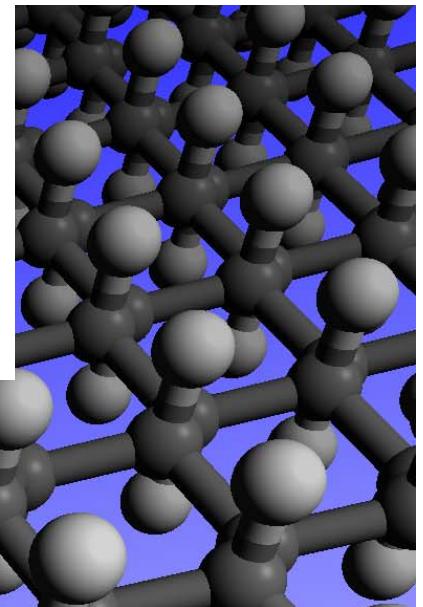
## Graphane: A two-dimensional hydrocarbon

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# Control of Graphene's Properties by Reversible Hydrogenation: Evidence for Graphane

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Although graphite is known as one of the most chemically inert materials, we have found that graphene, a single atomic plane of graphite, can react with atomic hydrogen, which transforms this highly conductive zero-overlap semimetal into an insulator. Transmission electron microscopy reveals that the obtained graphene derivative (graphane) is crystalline and retains the hexagonal lattice, but its period becomes markedly shorter than that of graphene. The reaction with hydrogen is reversible, so that the original metallic state, the lattice spacing, and even the quantum Hall effect can be restored by annealing. Our work illustrates the concept of graphene as a robust atomic-scale scaffold on the basis of which new two-dimensional crystals with designed electronic and other properties can be created by attaching other atoms and molecules.

Graphene, a flat, tightly packed layer of carbon atoms, continues to attract attention because of its unique properties and effects that arise from its fine structure (1). Chemical modifications of graphene have been less explored, even though carbon nanotubes suggest that they can be altered chemically without breaking the carbon bonds. For example, graphene oxide is graphene with its surface densely covered with hydroxyl and other groups (2–6). Unfortunately, graphene oxide is strongly disordered, poorly conductive, and difficult to

material. Particularly elegant is the idea of attaching atomic hydrogen to each site of the graphene lattice to create graphane (7), which changes the hybridization of carbon atoms from  $sp^2$  into  $sp^3$ ,

thus removing the conducting  $\pi$ -bands and opening an energy gap (7, 8).

Previously, absorption of hydrogen on graphitic surfaces was investigated mostly in con-

junction with hydrogen storage, with the research focused on physisorbed molecular hydrogen (9–11). More recently, atomic hydrogen chemisorbed on carbon nanotubes has been studied theoretically (12) as well as by a variety of experimental techniques including infrared (13), ultraviolet (14, 15), and x-ray (16) spectroscopy and scanning tunneling microscopy (17). We report the reversible hydrogenation of single-layer graphene and observed dramatic changes in its transport properties and in its electronic and atomic structure, as evidenced by Raman spectroscopy and transmission electron microscopy (TEM).

Graphene crystals were prepared by use of micromechanical cleavage (18) of graphite on top of an oxidized Si substrate (300 nm  $SiO_2$ ) and then identified by their optical contrast (1, 18) and distinctive Raman signatures (19). Three types of samples were used: large ( $>20 \mu m$ ) crystals for Raman studies, the standard Hall bar devices 1  $\mu m$  in width (18), and free-standing mem-

branes for TEM. For details of sample preparation see earlier work (18, 20, 21).

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0FA, UK. <sup>5</sup>Institute for Molecules and Materials, Radboud University Nijmegen, 6525 ED Nijmegen, Netherlands.

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# Control of Graphene's Properties by Reversible Hydrogenation: Evidence for Graphane

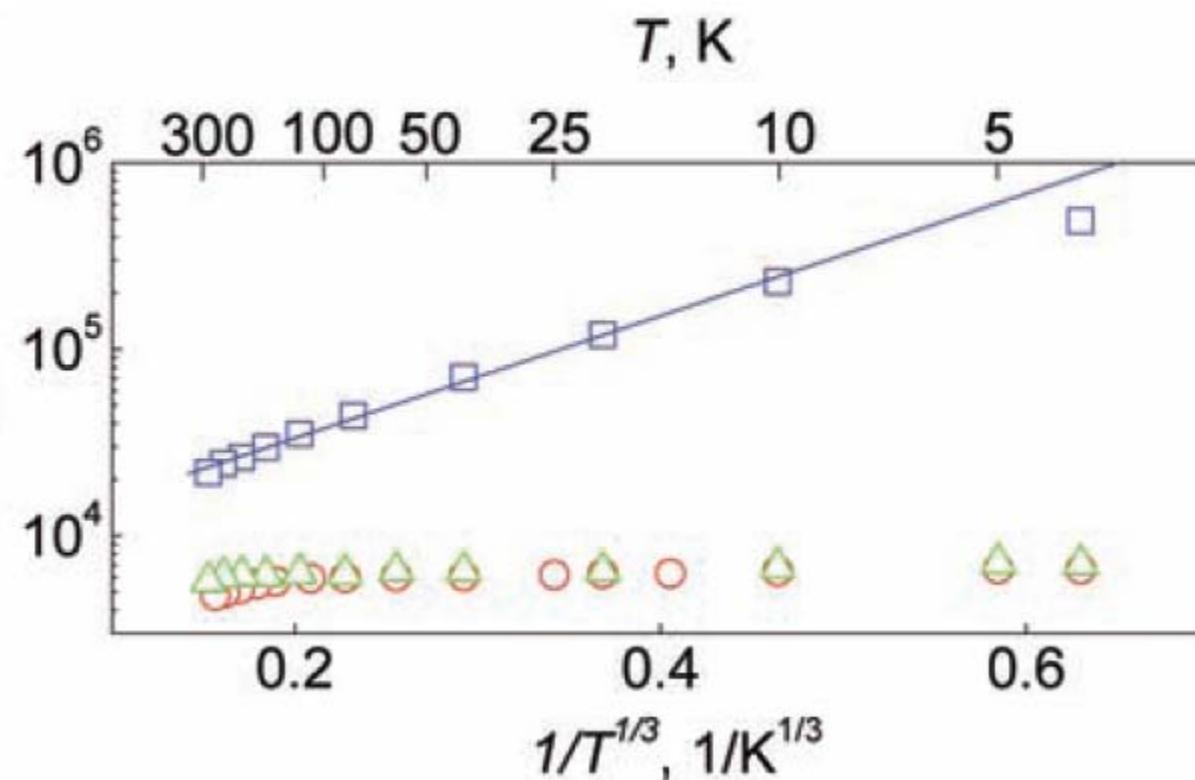
D. C. Elias,<sup>1,\*</sup> R. R. Nair,<sup>1,\*</sup> T. M. G. Mohiuddin,<sup>1</sup> S. V. Morozov,<sup>2</sup> P. Blake,<sup>3</sup> M. P. Halsall,<sup>1</sup> A. C. Ferrari,<sup>4</sup> D. W. Boukhvalov,<sup>5</sup> M. I. Katsnelson,<sup>5</sup> A. K. Geim,<sup>1,3</sup> K. S. Novoselov<sup>1†</sup>

Although graphite is known as one of the most chemically inert materials, we have found that graphene, a

highly conductive form of carbon, reveals that it is reversible, and its effect can be observed at the atomic-scale and other pr

**G**raphene has a very tight binding between carbon atoms because of the strong covalent bonds that form the hexagonal lattice. This tight binding makes it a very good conductor of electricity. However, it has been less explored than other forms of carbon nanotubes due to its relatively low density of states.

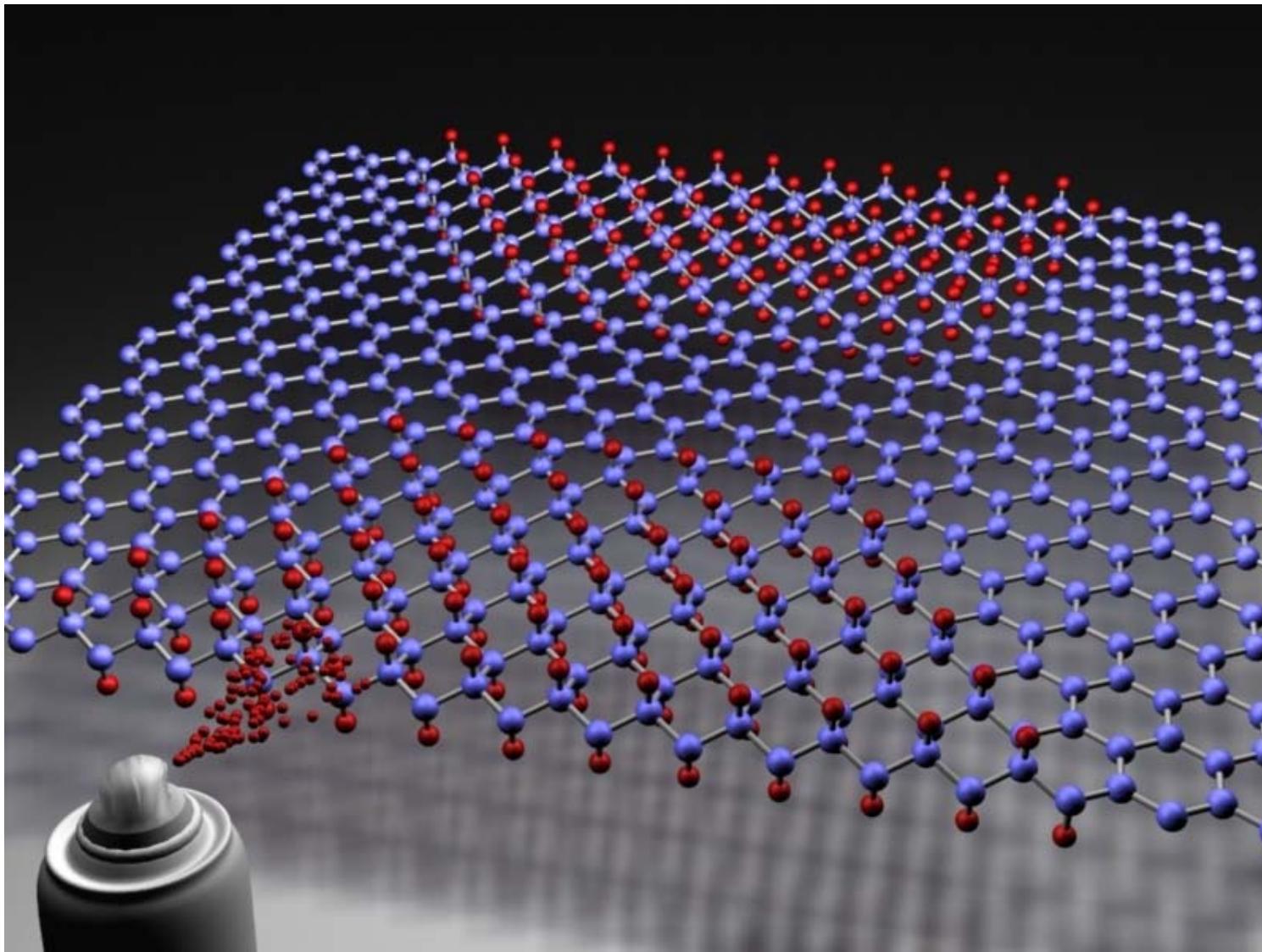
**C**onsequently, the properties of graphene are still not fully understood.



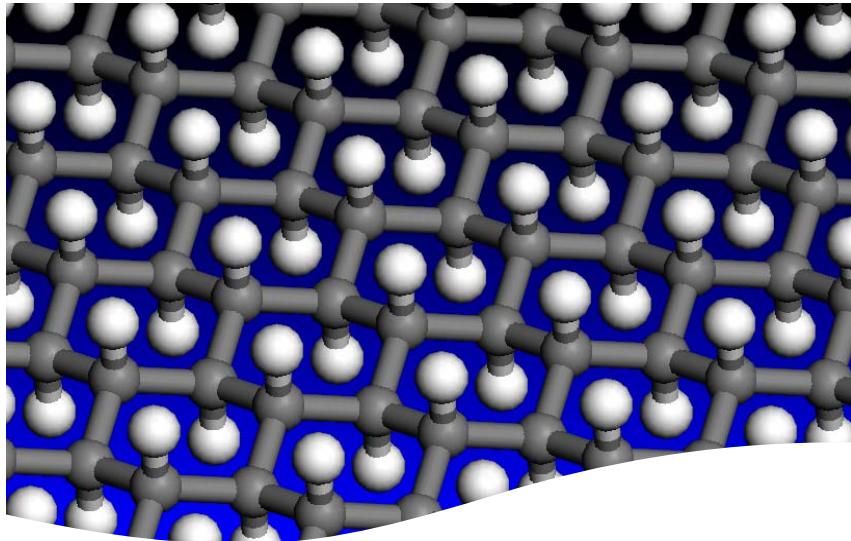
junction with hydrogen storage, with the research focused on physisorbed molecular hydrogen (9–11). More recently, atomic hydrogen chemisorbed on carbon nanotubes has been studied theoretically (12) as well as by a variety of experimental techniques including infrared (13), ultraviolet (14, 15), and x-ray (16) spectroscopy and scanning tunneling microscopy (17). We report the reversible hydrogenation of single-layer graphene and observed dramatic changes in its transport properties and in its electronic and atomic structure as evidenced by Raman spectroscopy and M).

by use of graphite on  $\text{SiO}_2$ ) and ast (1, 18) Three types n) crystals all bar dede of sample '8, 20, 21).

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Graphics courtesy of  
Kostya Novoselov



## CH → Graphane

Predicted by

JOS, A. Chaudhari, and G. Barber PRB **75**, 153401 (2007)

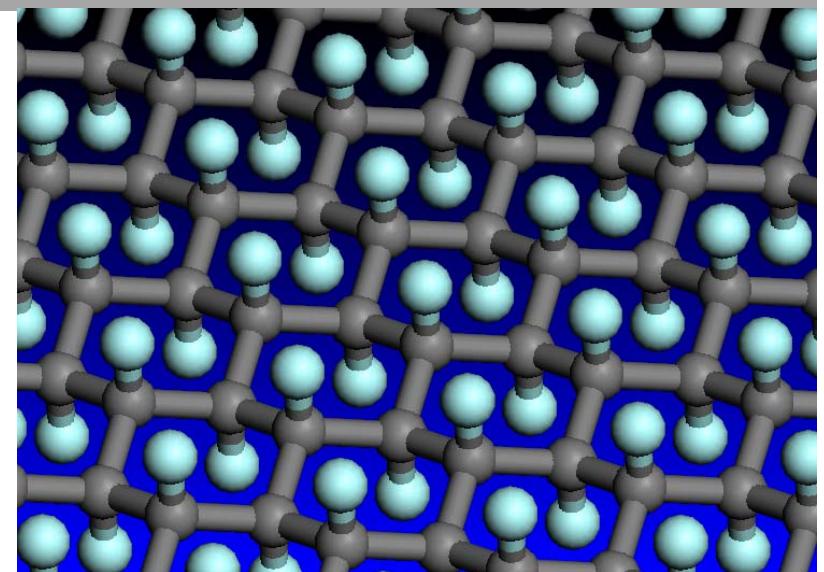
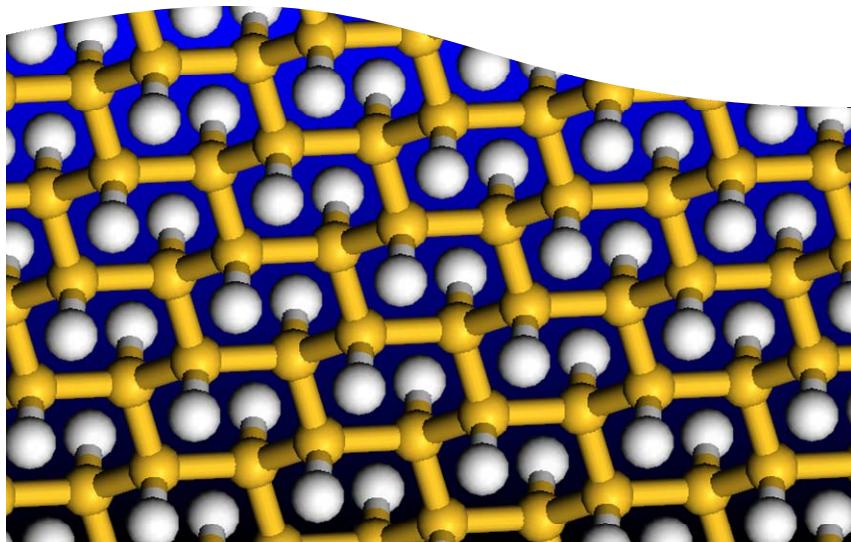
Some evidence shown by

D. C. Elias, et al. Science **323**, 610 (2009)

## CF → Carbon Monofluoride

Discovered by

O. Ruff and O. Bretshneider [JOS] Org. Allg. Chem. **217**, 1 (1934)



## SiH → Polysilane

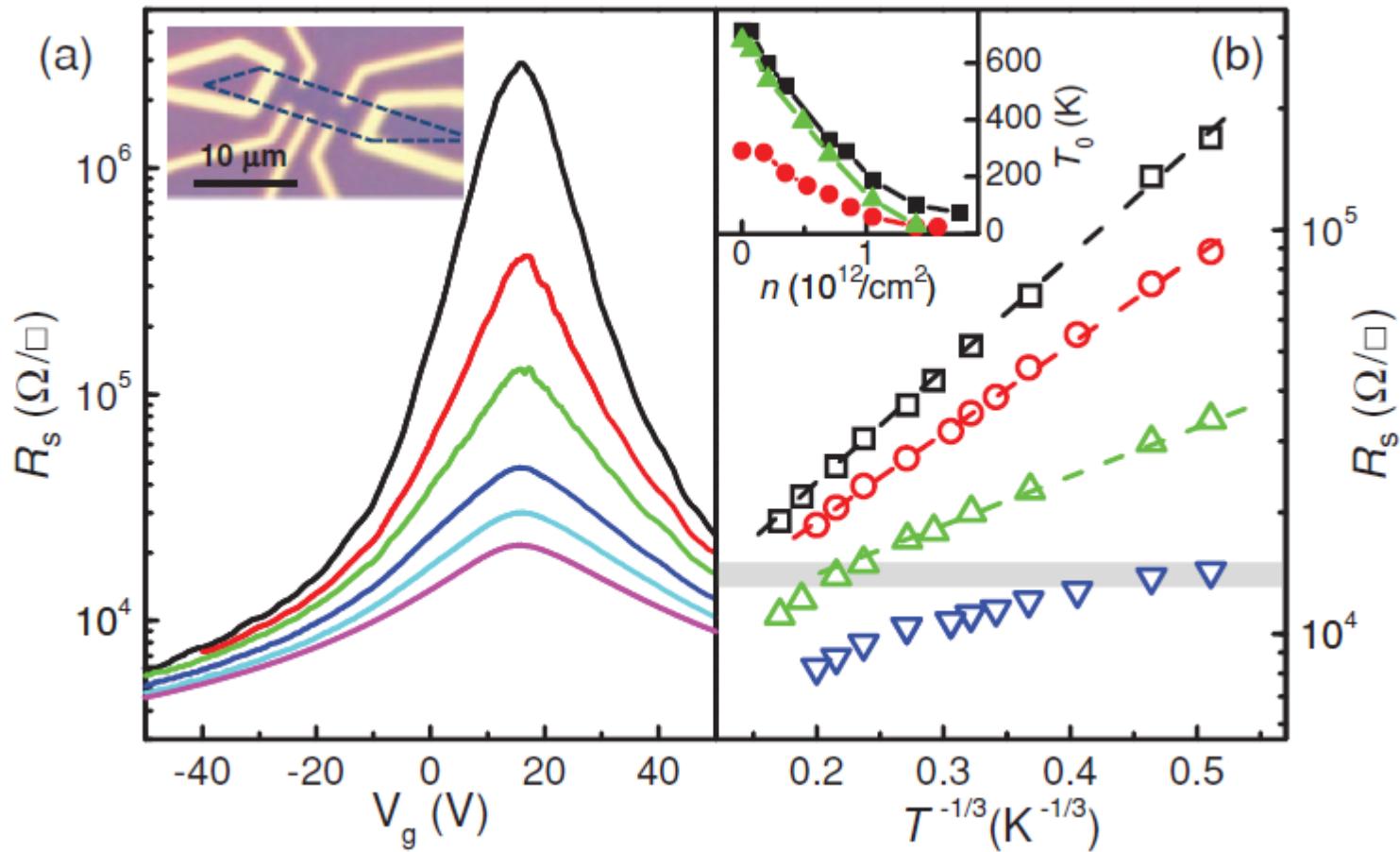
Predicted by

K. Takeda and K. Shiraishi, PRB **39**, 11028 (1989)

Synthesized by

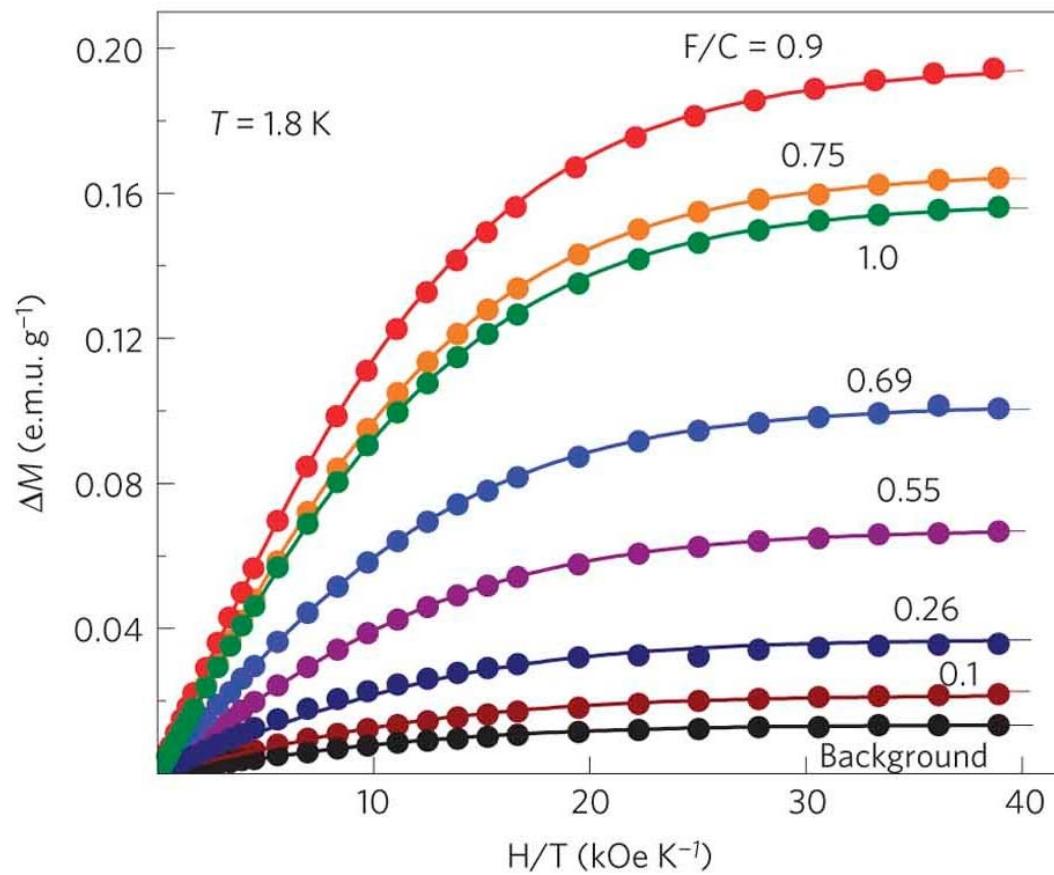
J. Dahn, B. Way, E. Fuller, and J. Tse PRB **48**, 17872 (1993)

# Partially fluorinated graphene



X. HONG, S.-H. CHENG, C. HERDING, AND J. ZHU

PHYSICAL REVIEW B 83, 085410 (2011)

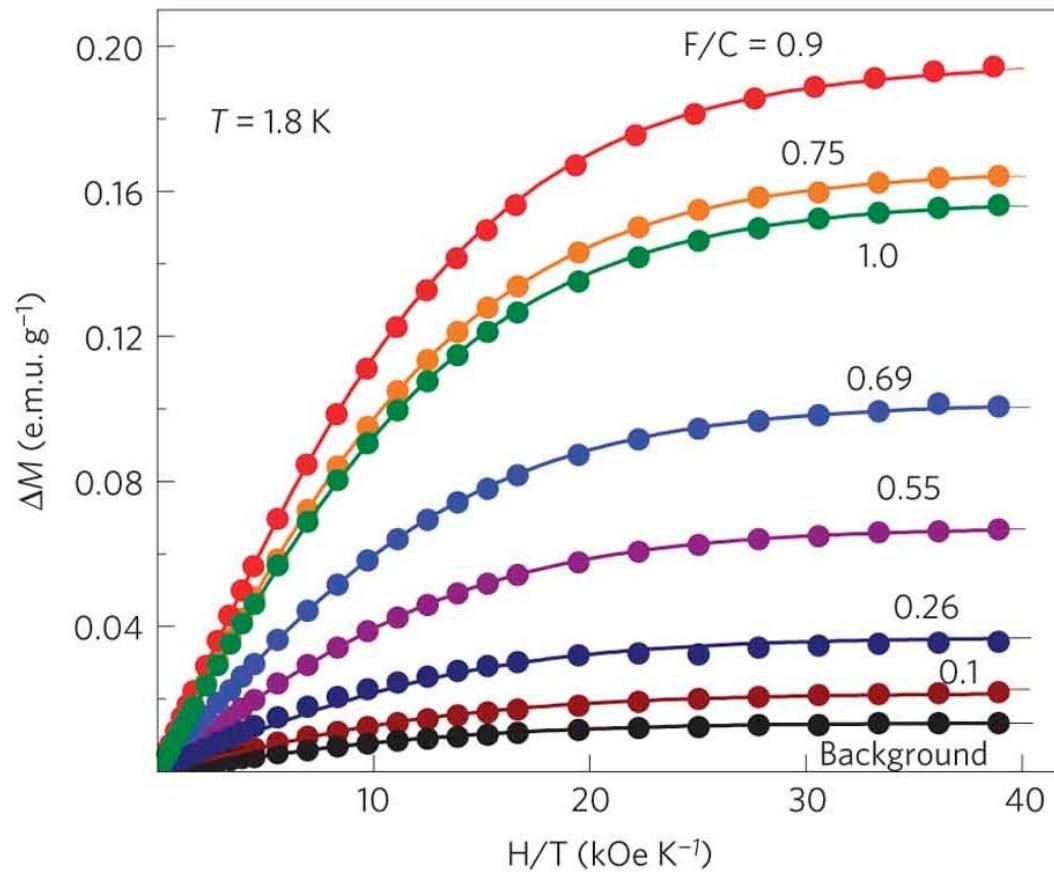


$$M = N g J \mu_B \left[ \frac{2J+1}{2J} \coth\left(\frac{2J+1}{2J} z\right) - \frac{1}{2J} \coth\left(\frac{z}{2J}\right) \right]$$

$$z = \frac{g J \mu_B H}{k_B T}$$

What about adsorbing more atoms...

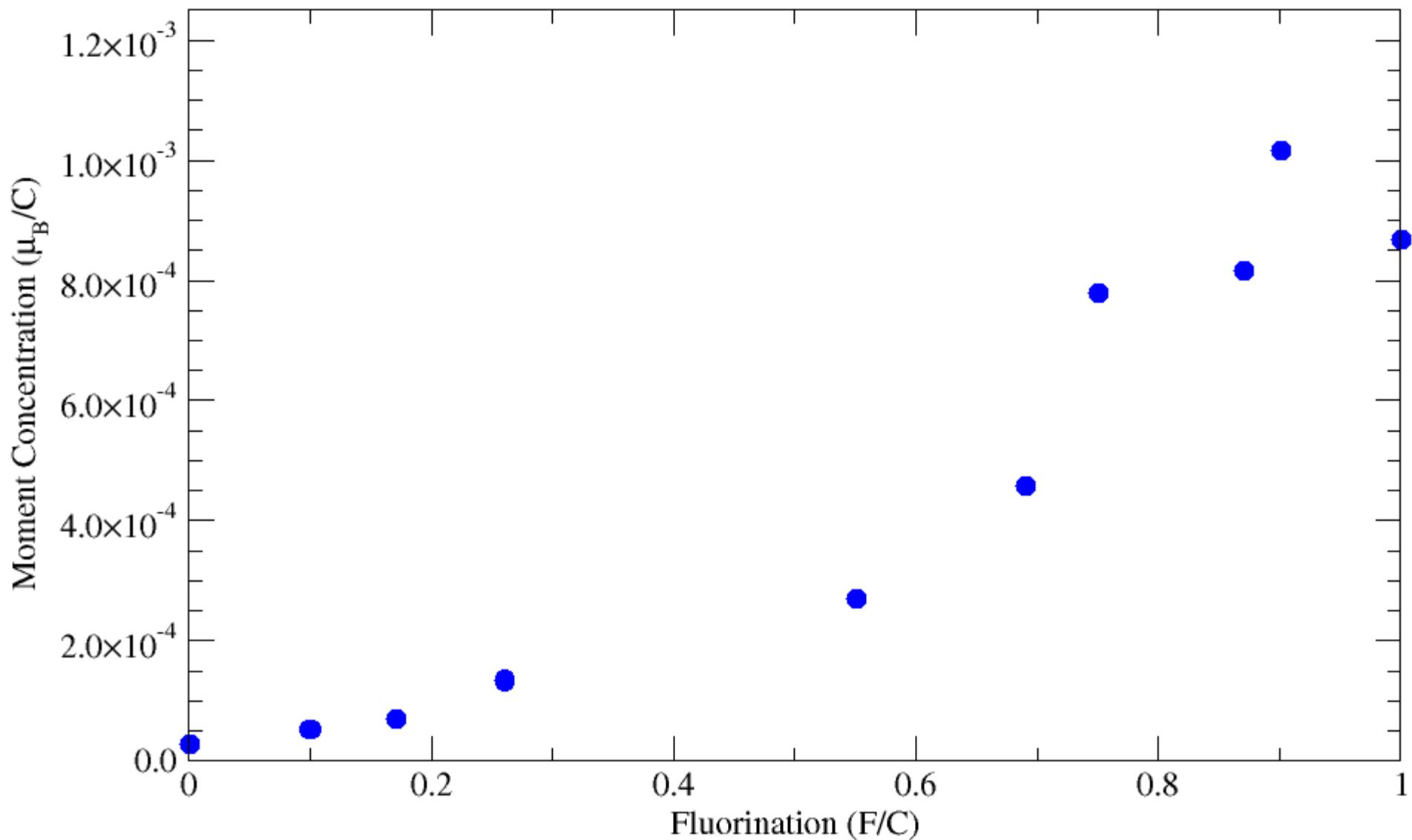
# IS FLUORINATED GRAPHENE MAGNETIC?

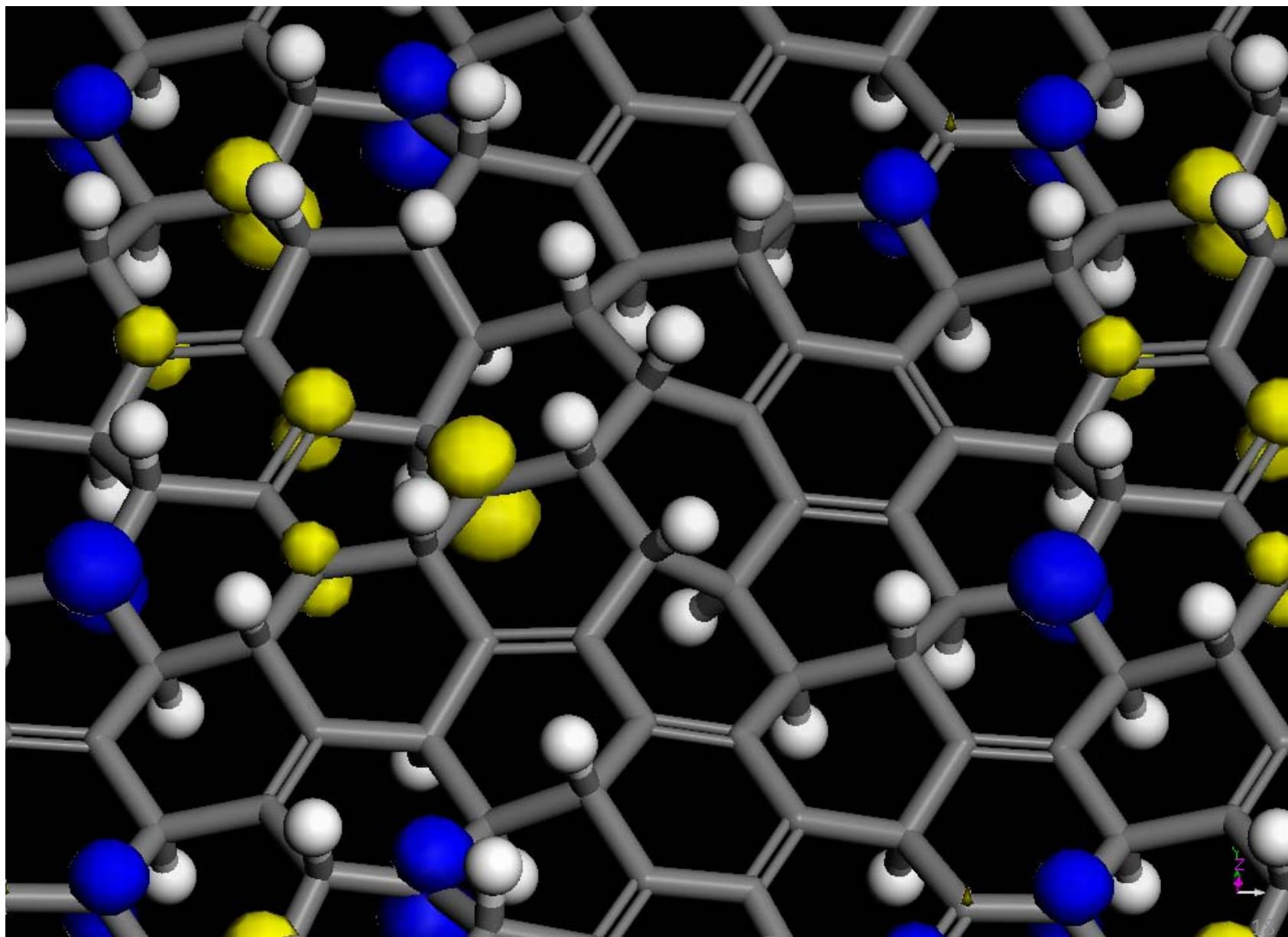


$$M = N g J \mu_B \left[ \frac{2J+1}{2J} \coth\left(\frac{2J+1}{2J} z\right) - \frac{1}{2J} \coth\left(\frac{z}{2J}\right) \right]$$

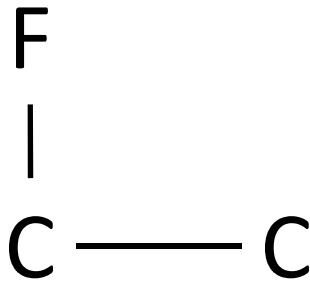
$$z = \frac{g J \mu_B H}{k_B T}$$

# Experimental fluorination sees magnetism...



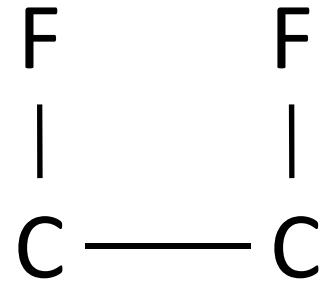


a)



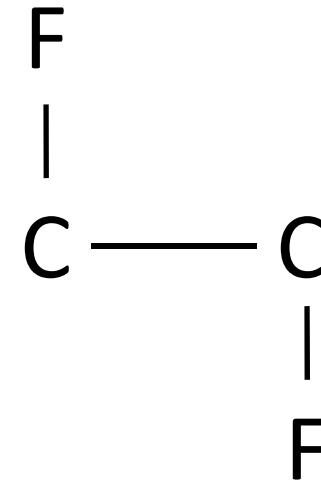
$$dE = 0$$

b)



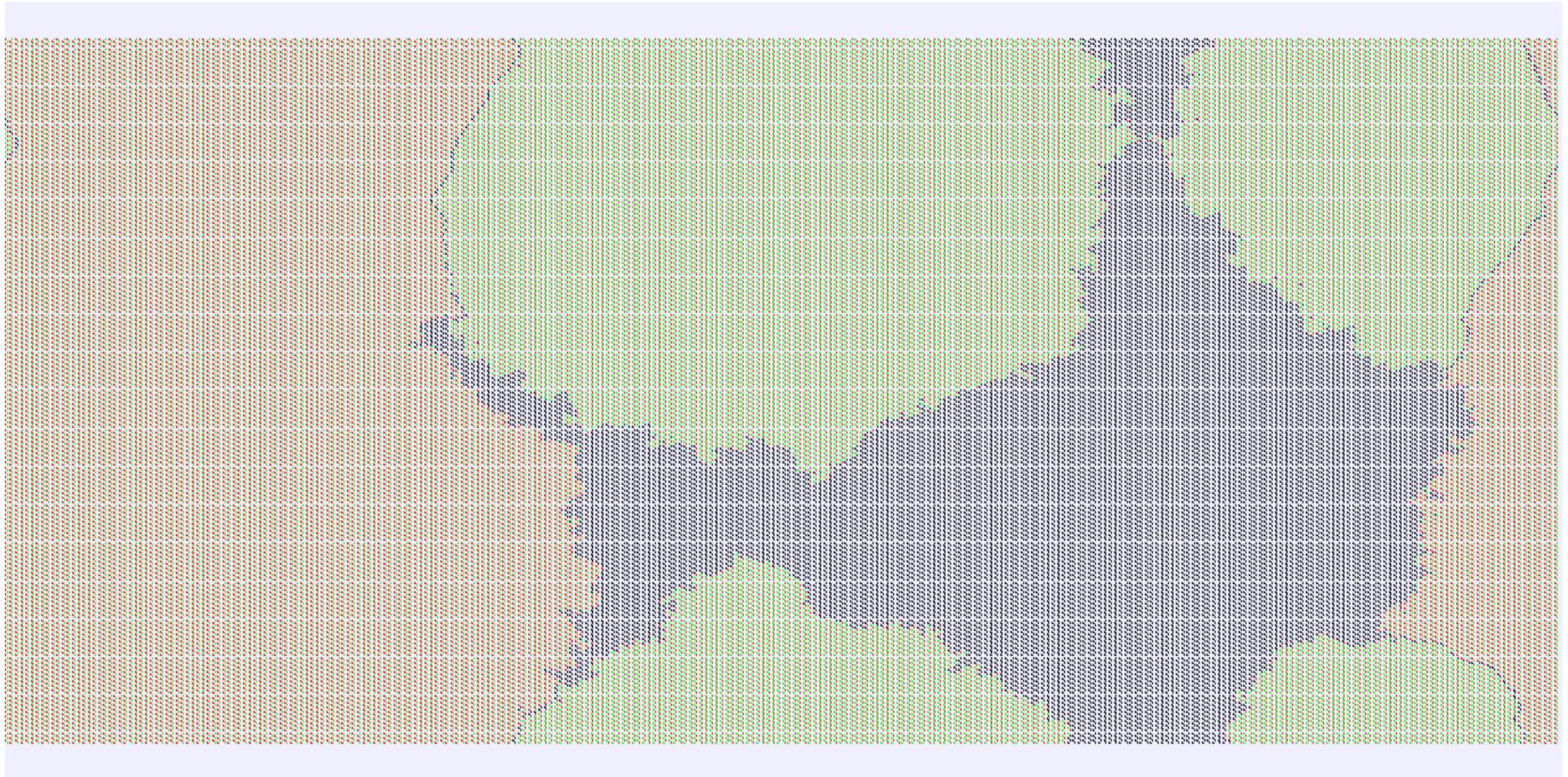
$$dE = J_{same}$$

c)

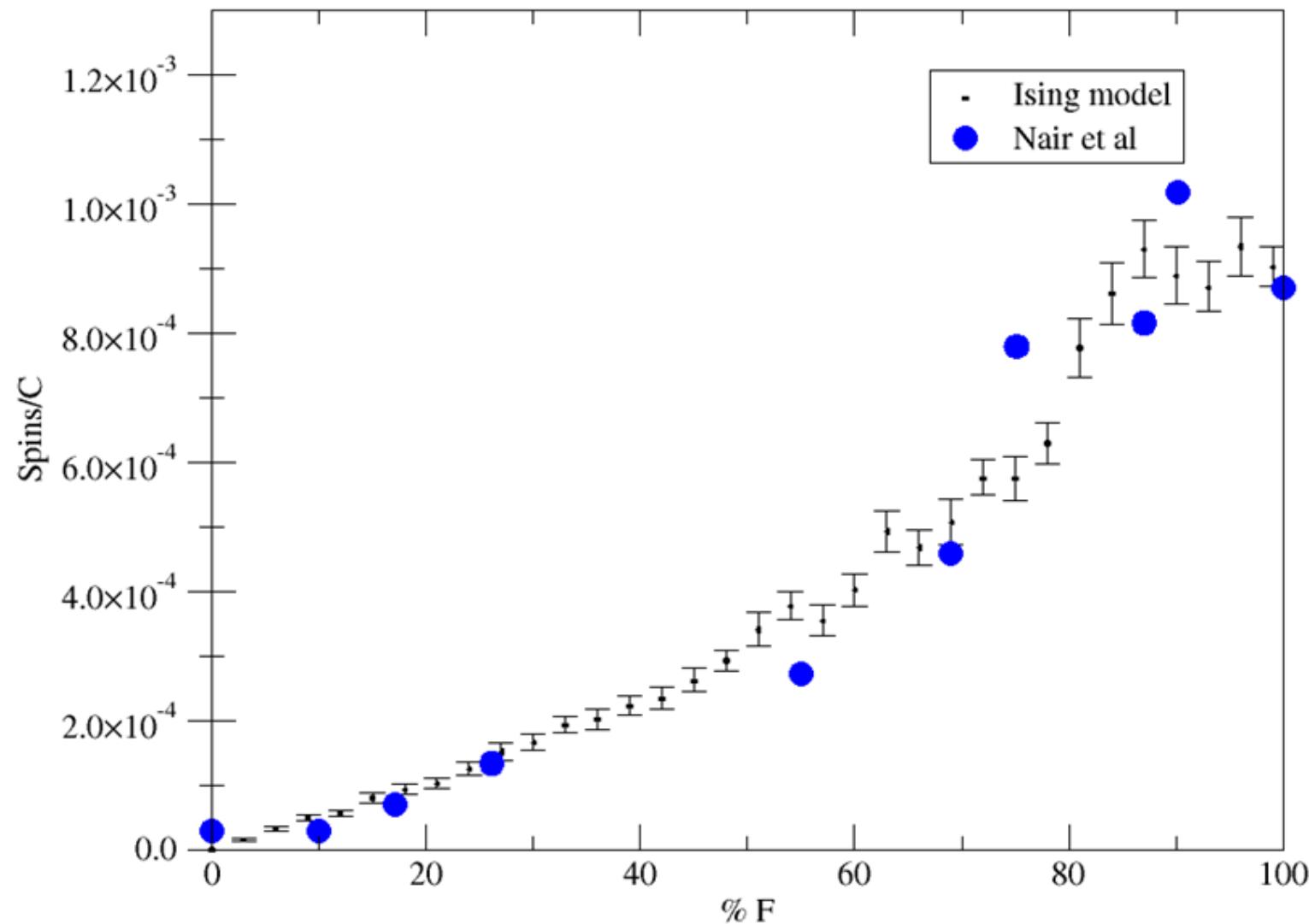


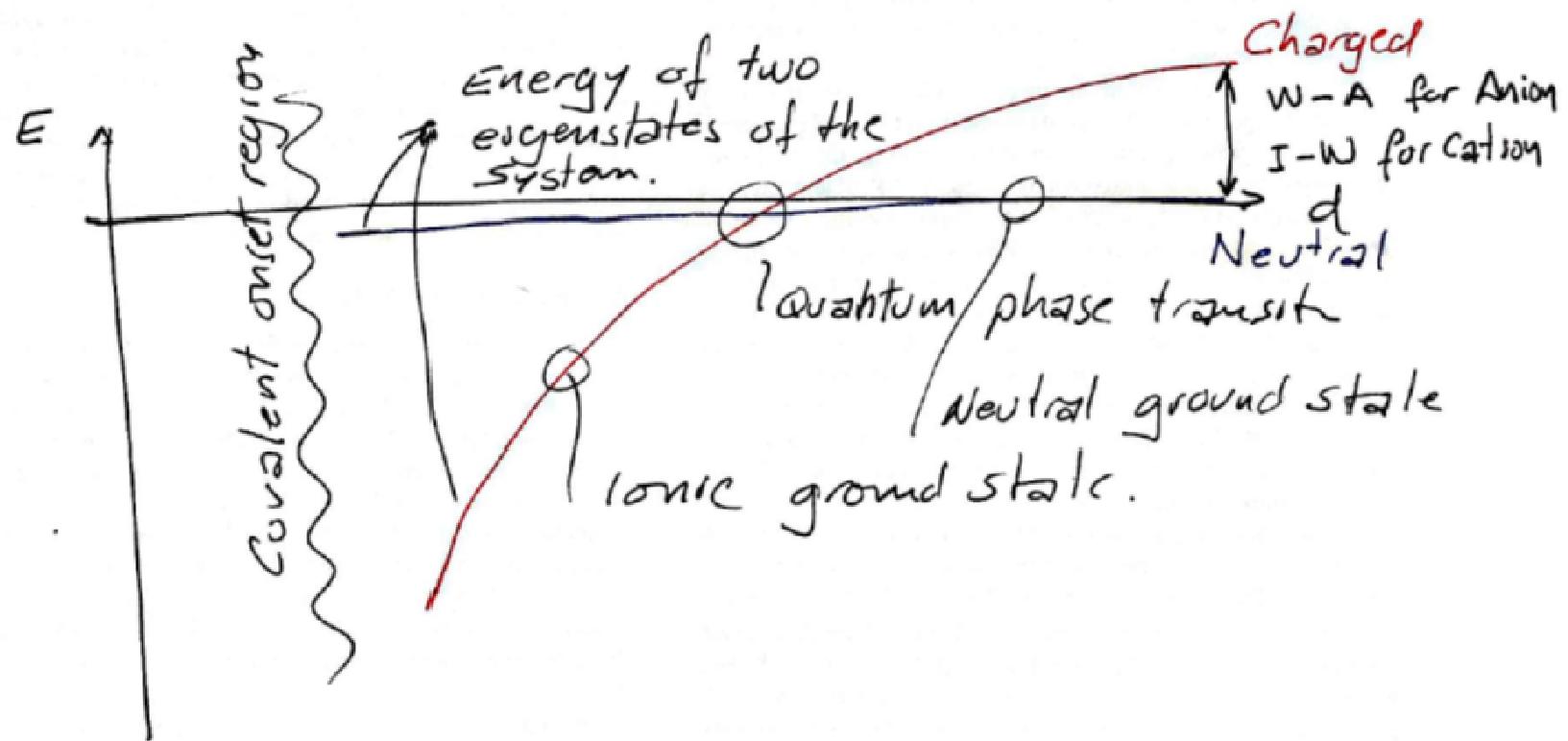
$$dE = J_{opp}$$

$$H^{IL} = \sum_{\langle i,j \rangle} (J_{opp} + (J_{same} - J_{opp})\delta_{s_i s_j}) s_i s_j - \mu \sum_i |s_i|$$



Ising Model  $J_{same} = 1, J_{opp} = -1, \mu = -1, T = 0.07$   
 $F/C = 0.75$





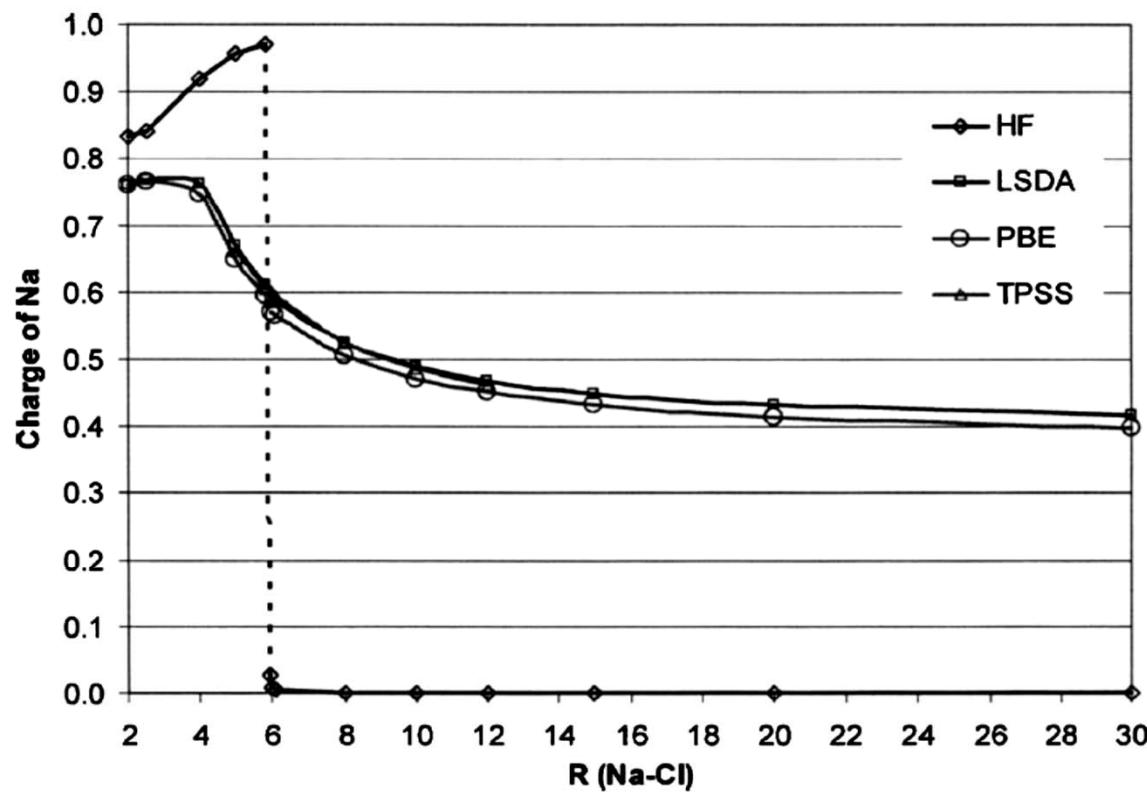
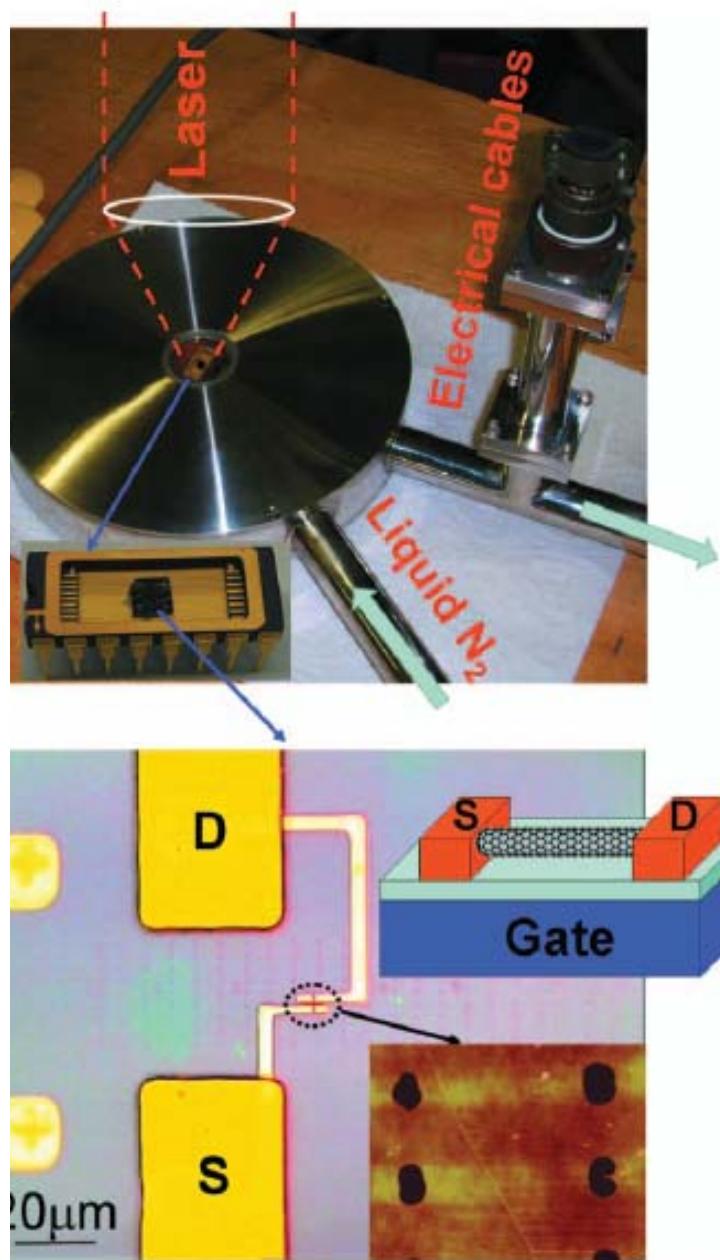


FIG. 3. Fractional (ChelpG) (Ref. 35) charge  $q$  on Na in NaCl as a function of bond length  $R$  (in Å). The Hartree-Fock “quantum phase transition” is qualitatively right, as discussed in the text.

Ruzsinszky, J. P. Perdew, G. I. Csonka, O. A. Vydrov, and G. E. Scuseria. 2006.  
 “Spurious Fractional Charge on Dissociated Atoms:  
 Pervasive and Resilient Self-Interaction Error of Common Density Functionals.”  
*The Journal of Chemical Physics* 125 (19): 194112. doi:10.1063/1.2387954.

What about gas detection...

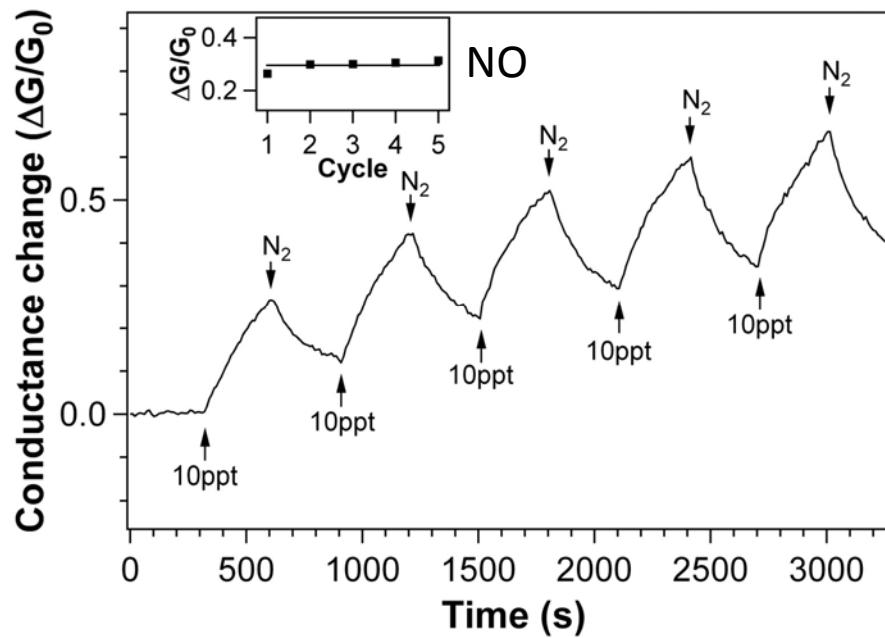
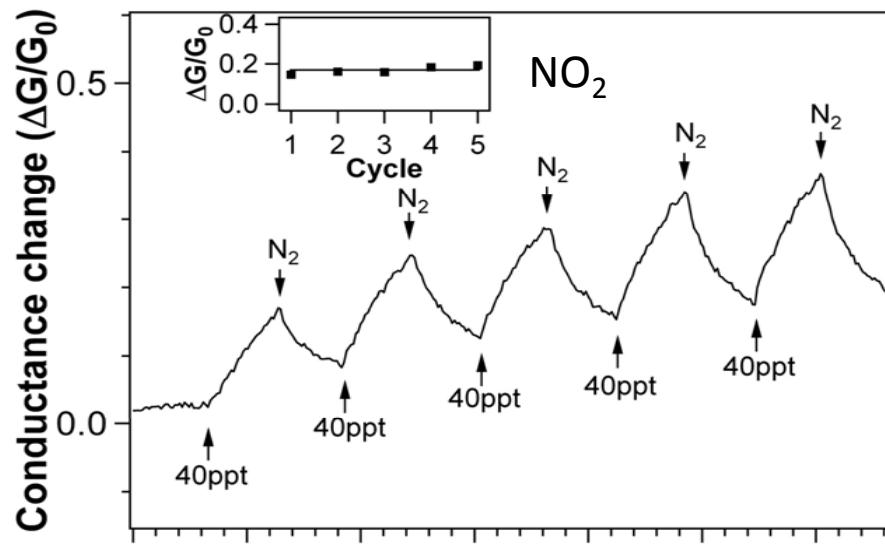
**VERY TINY AMOUNTS OF GAS DETECTED**

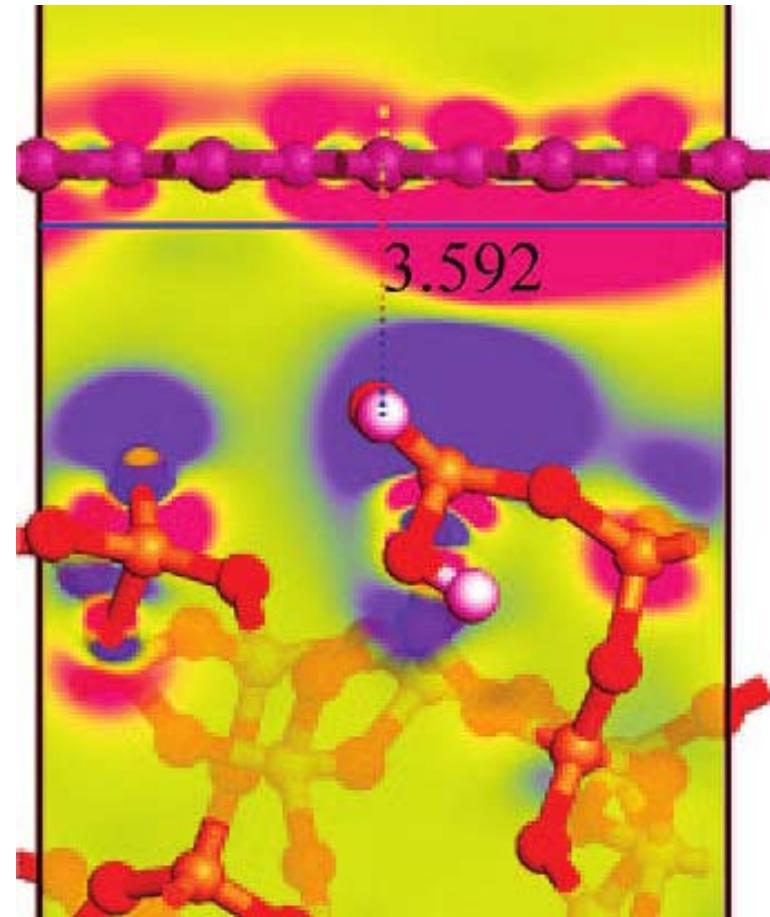
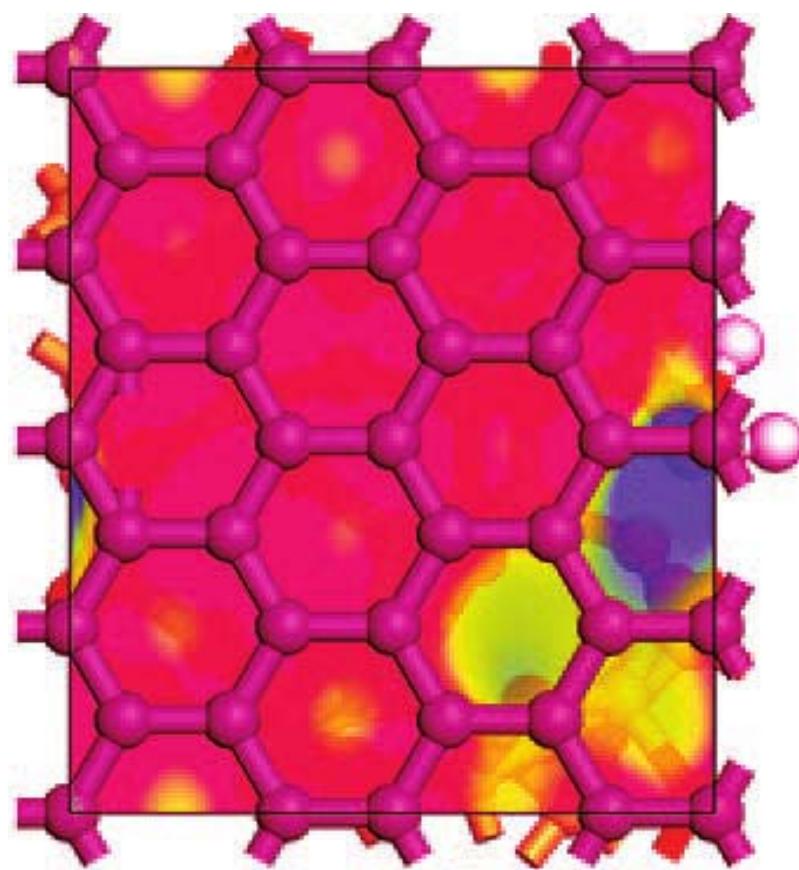


# HONDA

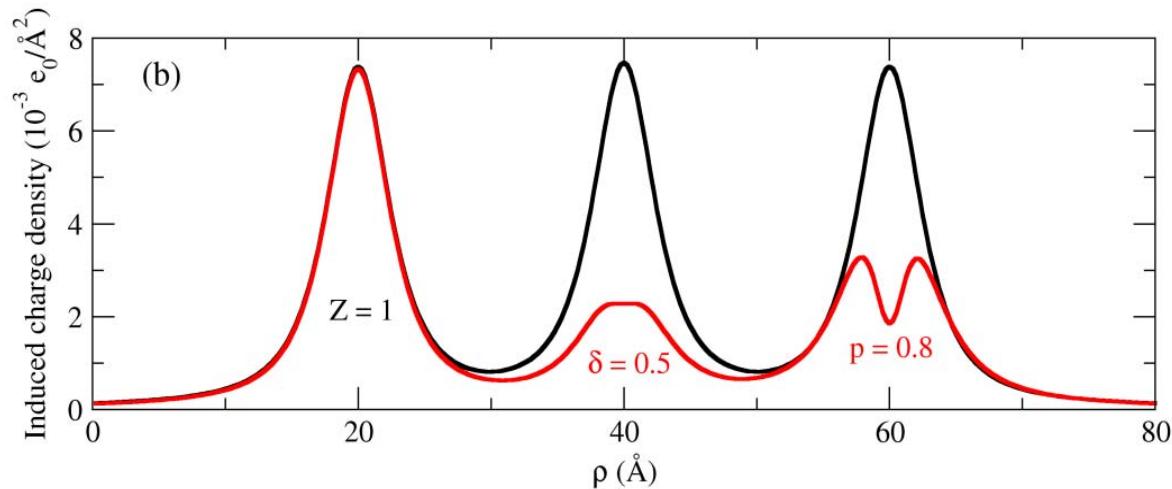
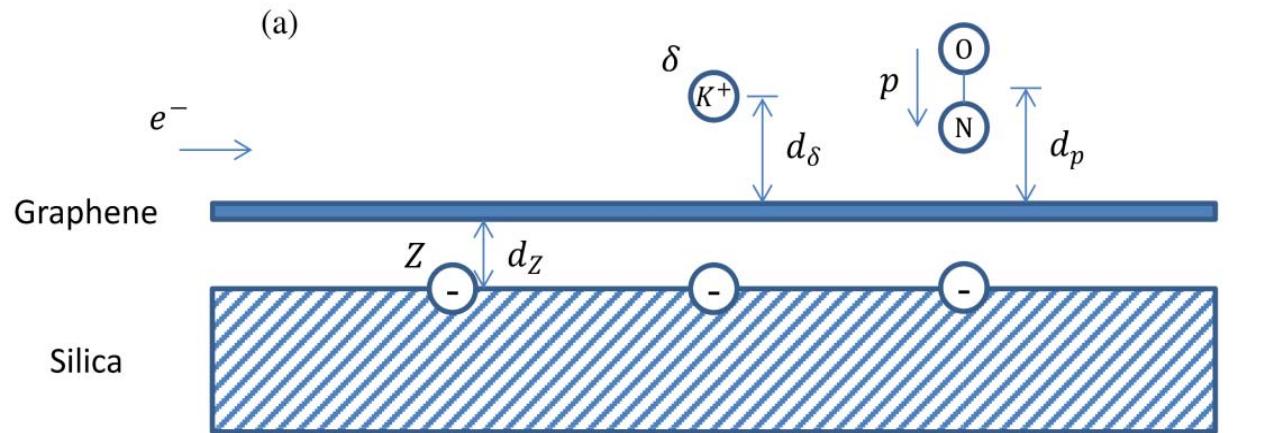
PENNSTATE  
1855

The Power of Dreams



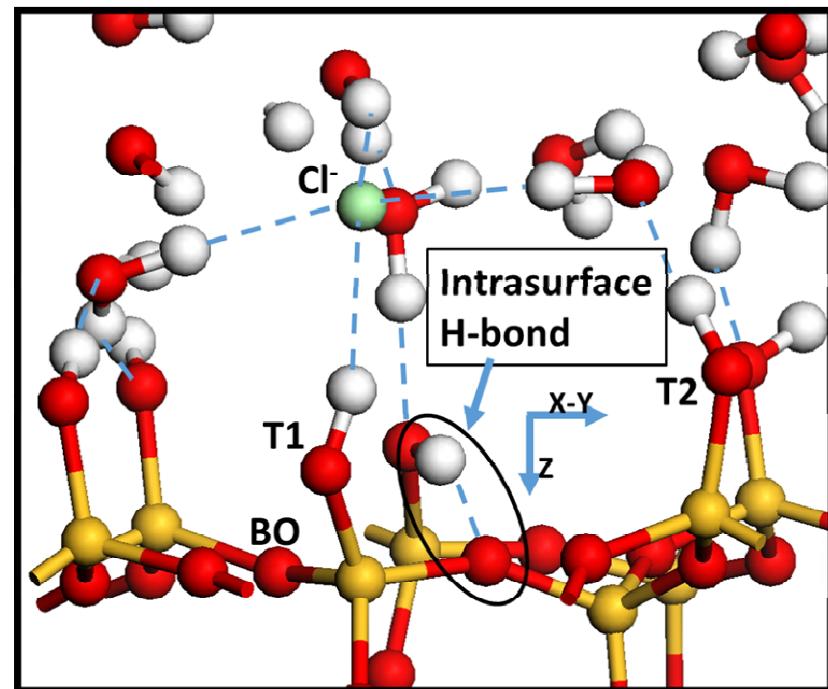
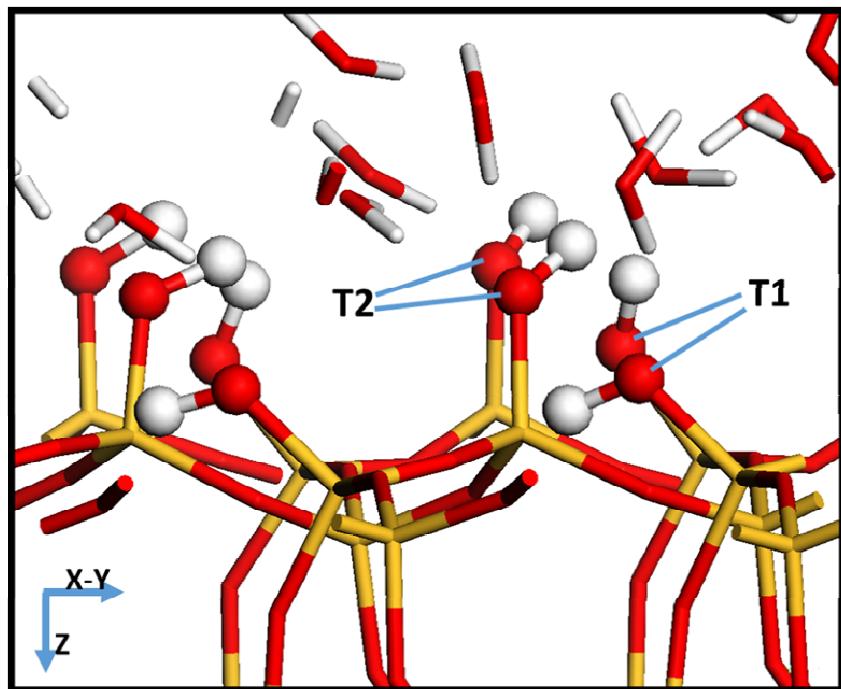


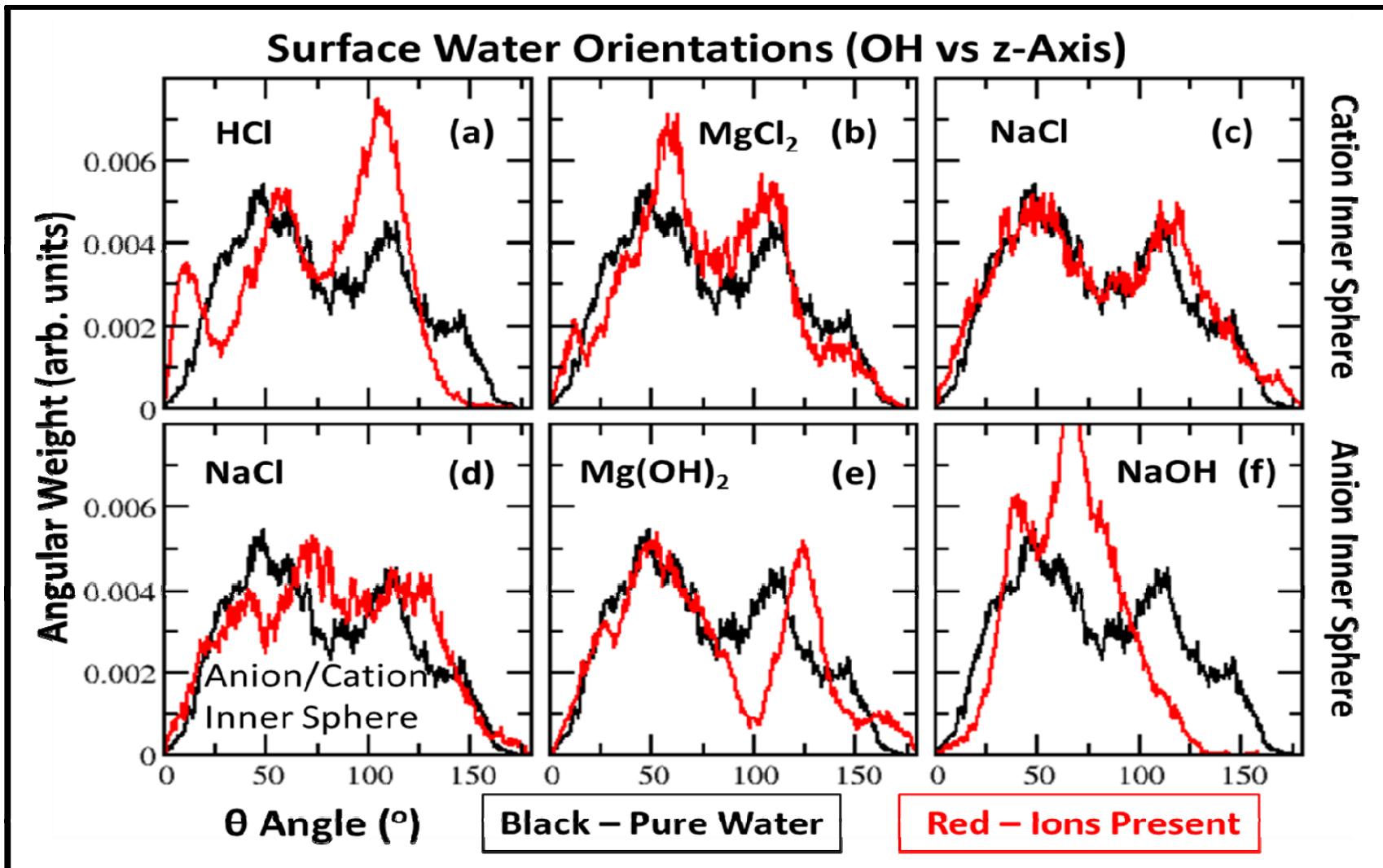
H. E. Romero, N. Shen, P. Joshi, H. R. Gutierrez, S. A. Tadigadapa, J. O. Sofo, and P. C. Eklund.  
“N-Type Behavior of Graphene Supported on Si/SiO<sub>2</sub> Substrates.” ACS Nano **2**, 2037 (2008)

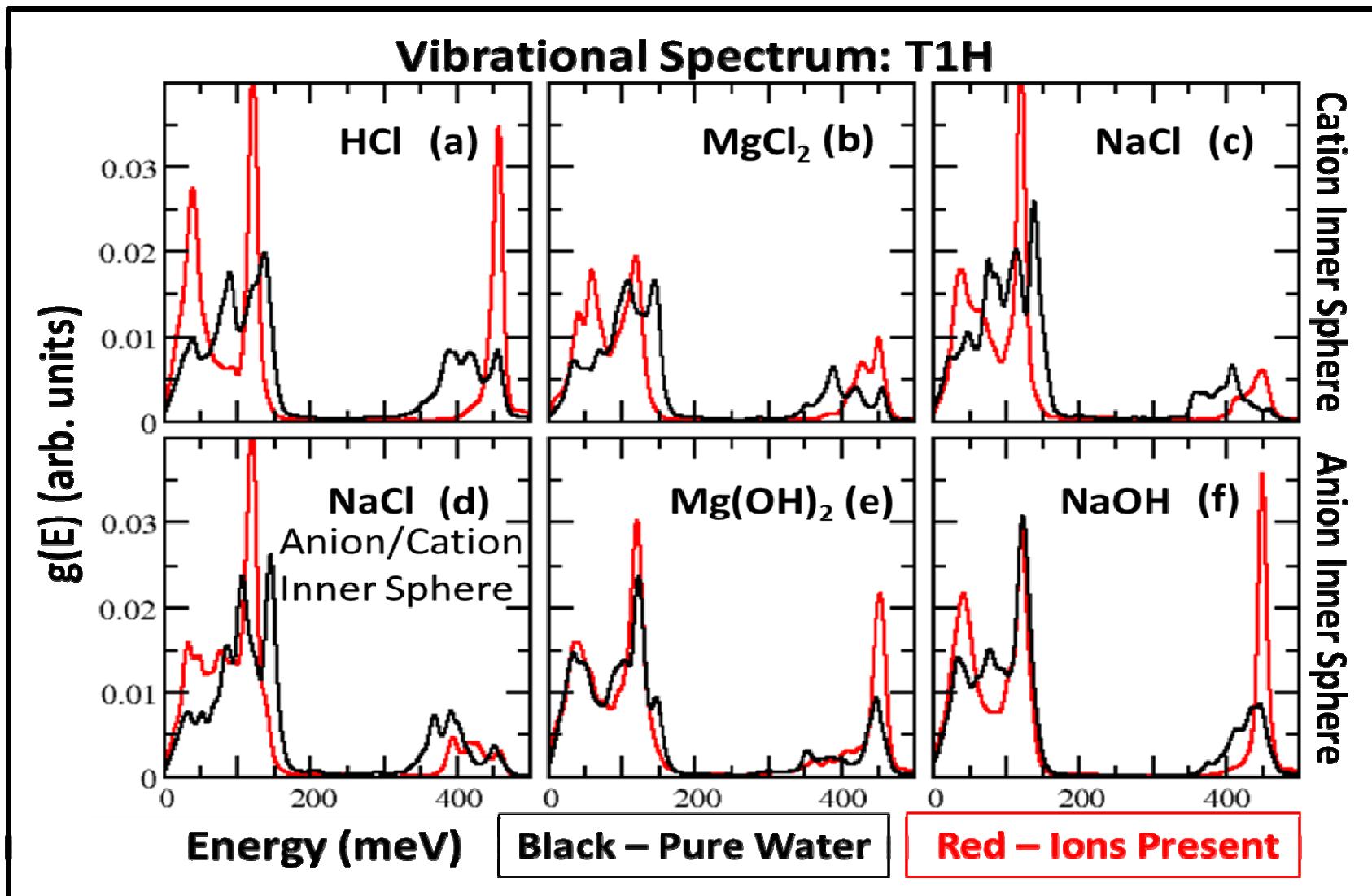


Silica dissolution...

# WHAT CAN WE SAY WITH ONE PICOSECOND

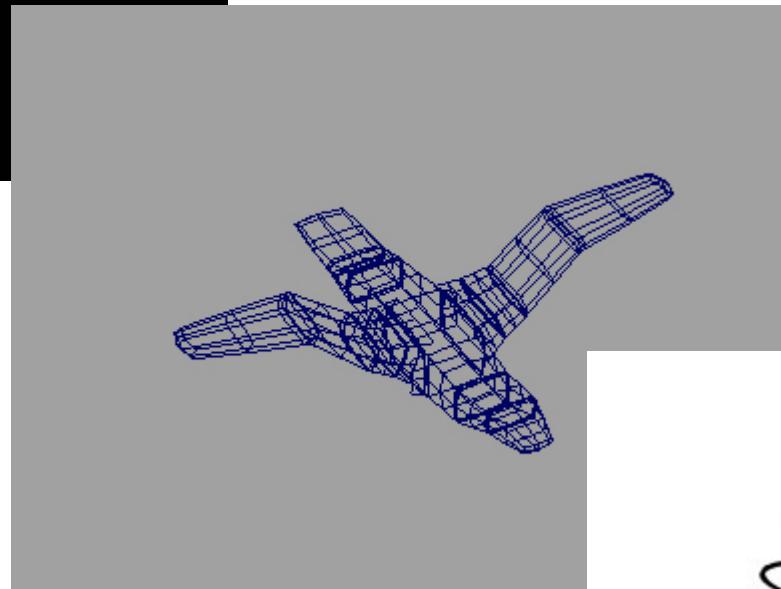






A final comment:

**EXPERIMENT<->SIMULATION<->THEORY**



Simulations describe complexity. Our theoretical work is to make it simple.